

# Performance Evaluation of F-barES-FEM-T4 in Dynamic Analysis

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# Background

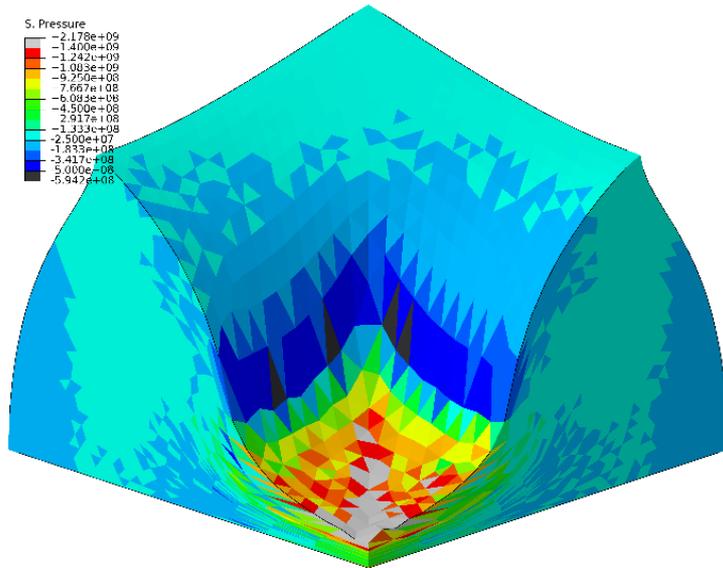
In the previous talk, F-bar aided edge-based smoothed finite element method with tetrahedral elements (**F-barES-FEM-T4**) is presented.

Characteristics of **F-barES-FEM-T4** in **static** analysis are as follows.

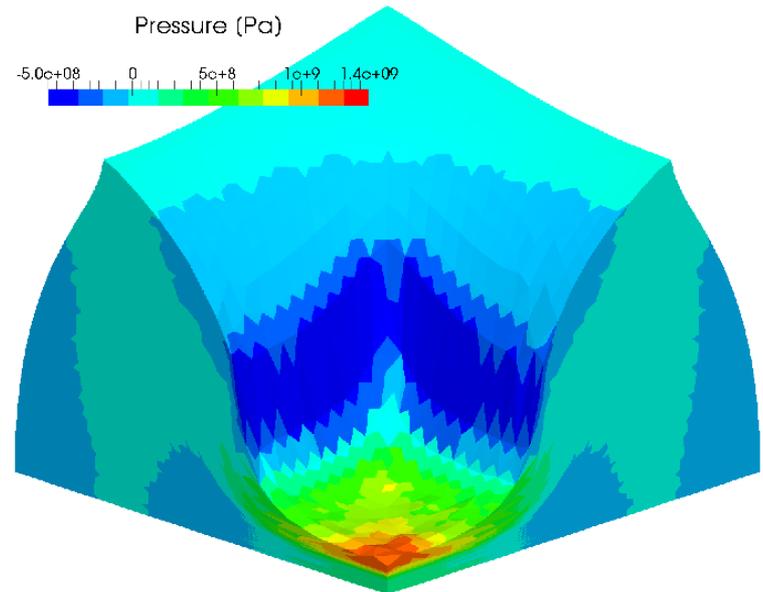
- ✓ Excellent accuracy,
- ✓ No increase in DOF,
- ✗ Increase in bandwidth of stiffness matrix.

# Characteristics (1 of 2)

✓ Excellent accuracy



ABAQUS C3D4H  
✗ pressure oscillation



F-barES-FEM-T4(3)  
# of cyclic smoothings

Our method shows excellent accuracy in **static** problems!

# Characteristics (2 of 2)

## ✓ No increase in DOF

- F-barES-FEM-T4 is a purely **displacement-based** formulation.
- In contrast to the u/p hybrid formulations, F-barES-FEM-T4 can be directly applied to explicit dynamics.

## ✗ Increase in bandwidth of stiffness matrix

- F-barES-FEM-T4 takes longer time to solve the equilibrium equations...
- In **explicit dynamics**, however, **we don't need to solve them!**

F-barES-FEM-T4 would be suitable  
for **explicit dynamics** of rubber-like materials!

# Objective

## Objective

Evaluate the performance of F-barES-FEM-T4  
in **explicit dynamics** for rubber-like materials.

## Table of Body Contents

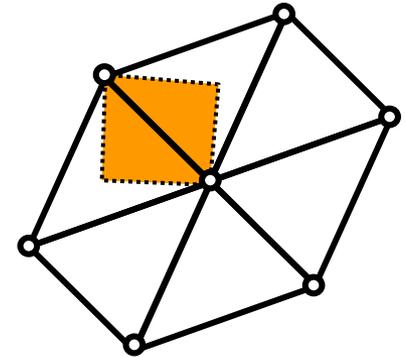
- Methods: Quick introduction of F-barES-FEM-T4
- Results & Discussion: A few verification analyses
- Summary

# Methods

# Procedure of F-bar ES-FEM (1 of 2)

Deformation gradient of **each edge**,  $\bar{\mathbf{F}}$  is derived as

$$\bar{\mathbf{F}} = \tilde{\mathbf{F}}^{\text{iso}} \cdot \bar{\mathbf{F}}^{\text{vol}}$$



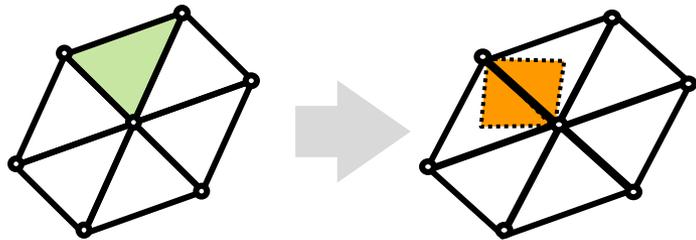
ES-FEM

# Procedure of F-bar ES-FEM (2 of 2)

Each part of  $\bar{F}$  is calculated as

$$\bar{F} = \tilde{F}^{\text{iso}} \cdot \bar{F}^{\text{vol}}$$

## Isovolumetric part

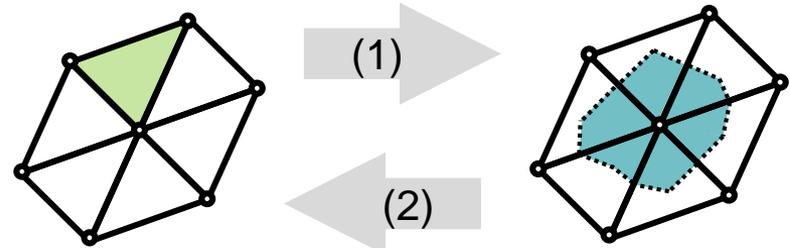


Smoothing the value of adjacent elements.



The same manner as  
ES-FEM

## Volumetric part



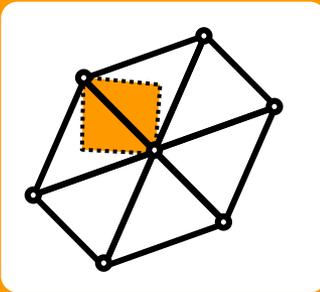
- (1) Calculating node's value by smoothing the value of adjacent elements
- (2) Calculating elements' value by smoothing the value of adjacent nodes
- (3) Repeating (1) and (2) a few times

# Advantages of F-bar ES-FEM

This formulation is designed to have 3 advantages.

$$\bar{\mathbf{F}} = \tilde{\mathbf{F}}^{\text{iso}} \cdot \bar{\mathbf{F}}^{\text{vol}}$$

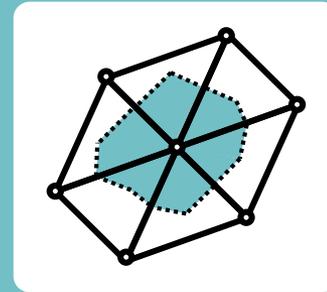
Isovolumetric part



Like a ES-FEM

1. Shear locking free

Volumetric part



Like a NS-FEM

2. Little pressure oscillation

3. Volumetric locking free  
with the aid of F-bar method

# Equation to solve

## Equation of Motion

$$[M]\{\ddot{\mathbf{u}}\} = \{\mathbf{f}^{\text{ext}}\} - \{\mathbf{f}^{\text{int}}\},$$

where

$$\{\mathbf{f}^{\text{int}}\} = \sum [\tilde{\mathbf{B}}]\{\bar{\mathbf{T}}\}V$$

B-matrix of ES-FEM

Stress derived from  $\bar{\mathbf{F}}$

Defenition of  $\{\mathbf{f}^{\text{int}}\}$   
in the same fashion  
as F-bar method

## Time integration

Velocity Verlet method (2<sup>nd</sup> order symplectic integrator)

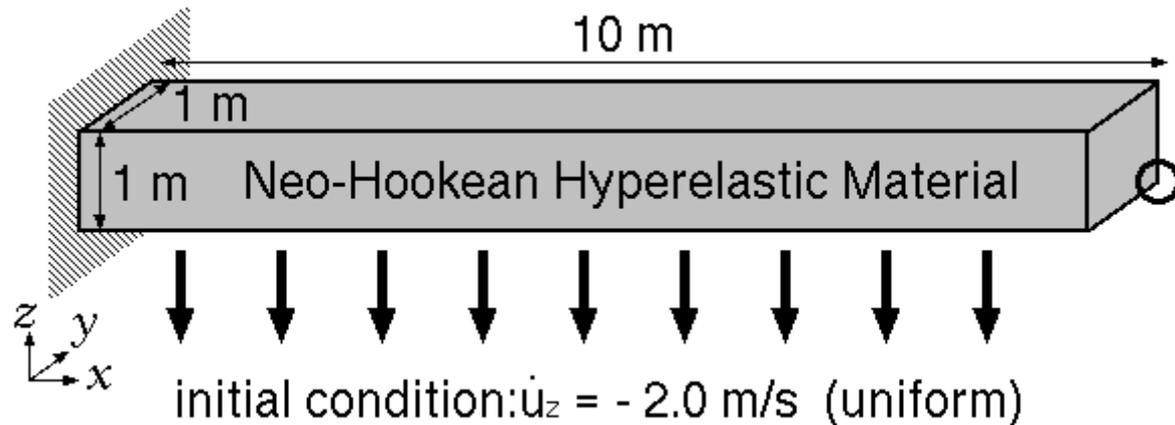
$$\{\mathbf{u}_{n+1}\} = \{\mathbf{u}_n\} + \{\dot{\mathbf{u}}_n\}\Delta t + \frac{1}{2}\{\ddot{\mathbf{u}}_n\}\Delta t^2$$

$$\{\ddot{\mathbf{u}}_{n+1}\} = [\mathbf{M}^{-1}](\{\mathbf{f}^{\text{ext}}\} - \{\mathbf{f}^{\text{int}}(\mathbf{u}_{n+1})\})$$

$$\{\dot{\mathbf{u}}_{n+1}\} = \{\dot{\mathbf{u}}_n\} + (\{\ddot{\mathbf{u}}_n\} + \{\ddot{\mathbf{u}}_{n+1}\})\Delta t/2$$

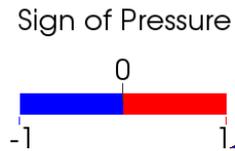
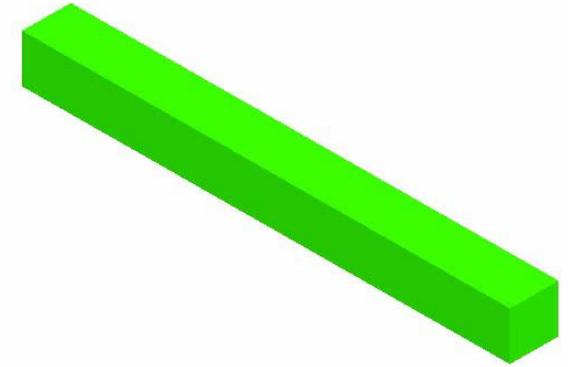
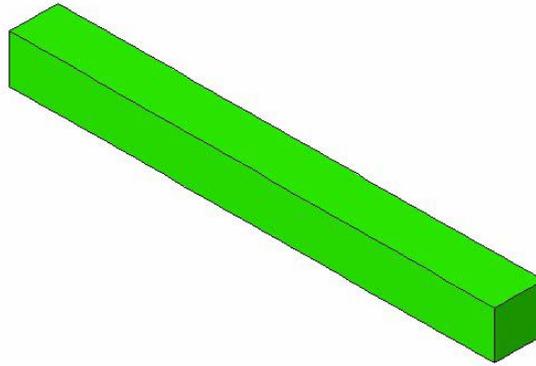
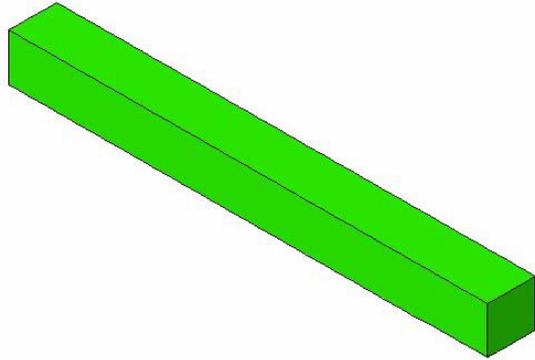
# Result & Discussion

# #1 Bending of a cantilever



- **Dynamic explicit** analysis.
- Neo-Hookean material
  - Initial Young's modulus: 6.0 MPa,
  - Initial Poisson's ratio: 0.499,
  - Density: 10000 kg/m<sup>3</sup>.
- Compare the results of F-barES-FEM-T4, Standard T4 (ABAQUS/Explicit C3D4) and Selective H8 (ABAQUS/Explicit C3D8) elements.

# Time history of deformed shapes



ABAQUS/Explicit C3D4  
(Standard T4 element)

- ✗ Pressure oscillation
- ✗ Locking

ABAQUS/Explicit C3D8  
(Selective H8 element)

**Reference**

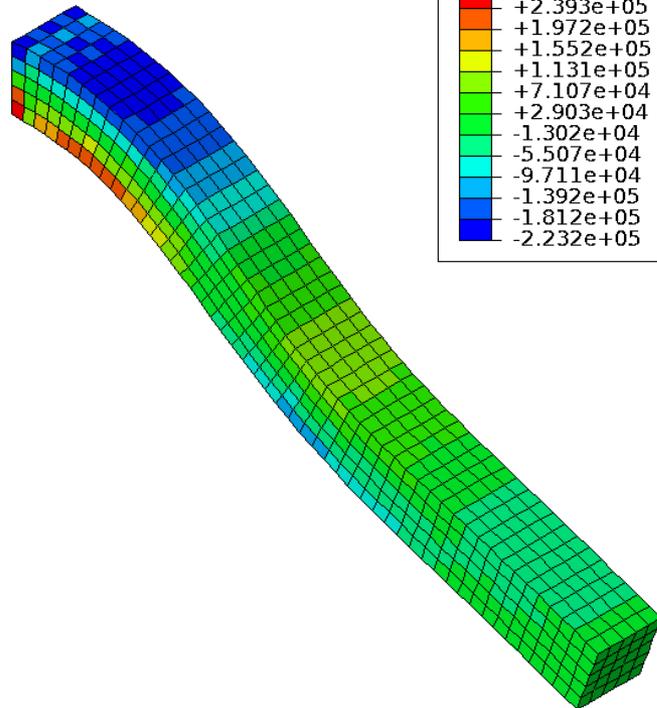
F-barES-FEM-T4(2)  
(Proposed method)

- ✓ No pressure oscillation
- ✓ No locking

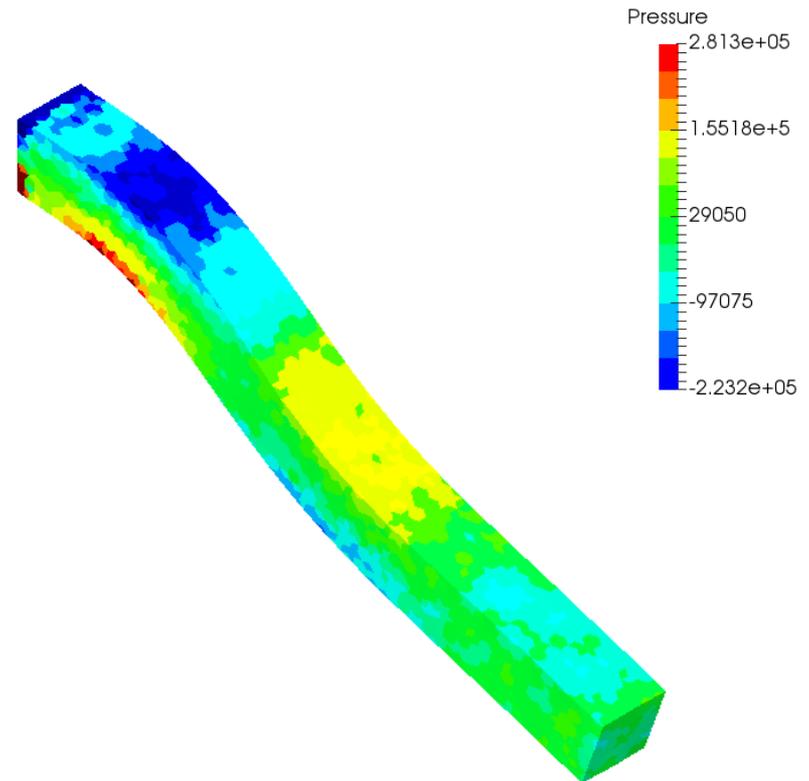
Proposed method suppresses pressure oscillation and locking!

# Deformed shapes and pressure distributions

at  $t = 1.5$  s



ABAQUS/Explicit C3D8  
(Selective H8 element)  
**Reference**

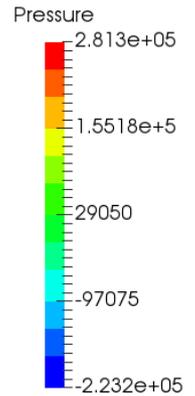
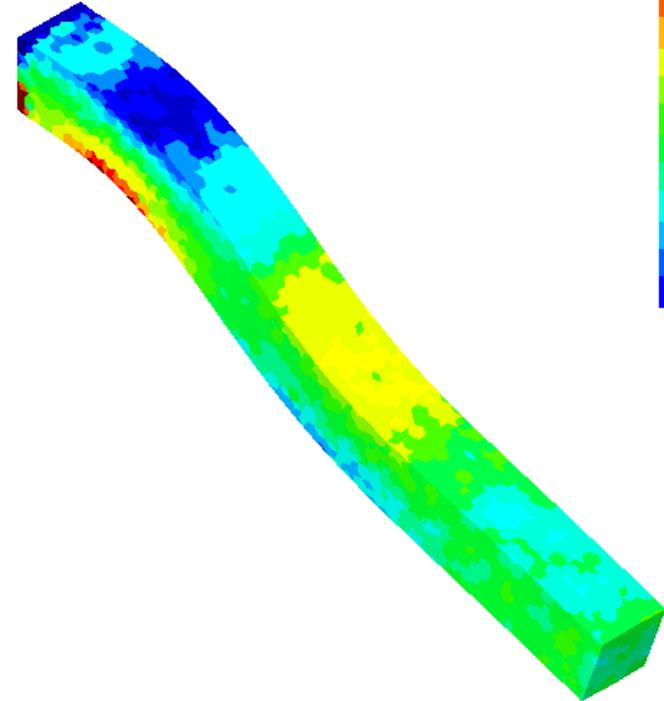
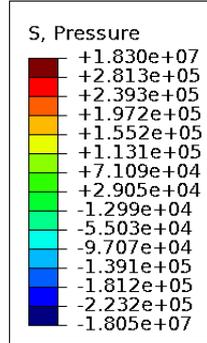
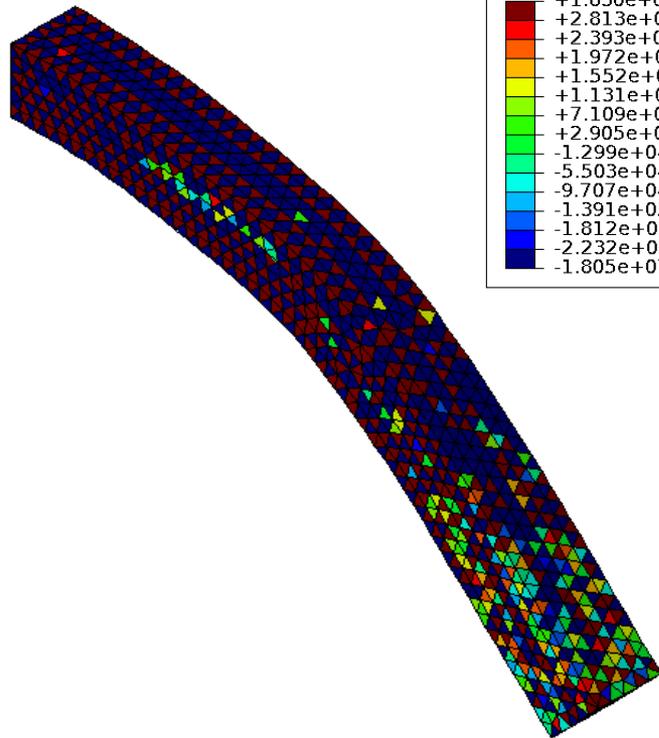


F-barES-FEM-T4(2)  
(Proposed method)

Proposed method is comparable to Selective H8 element!

# Deformed shapes and pressure distributions

at  $t = 1.5$  s



ABAQUS/Explicit C3D4

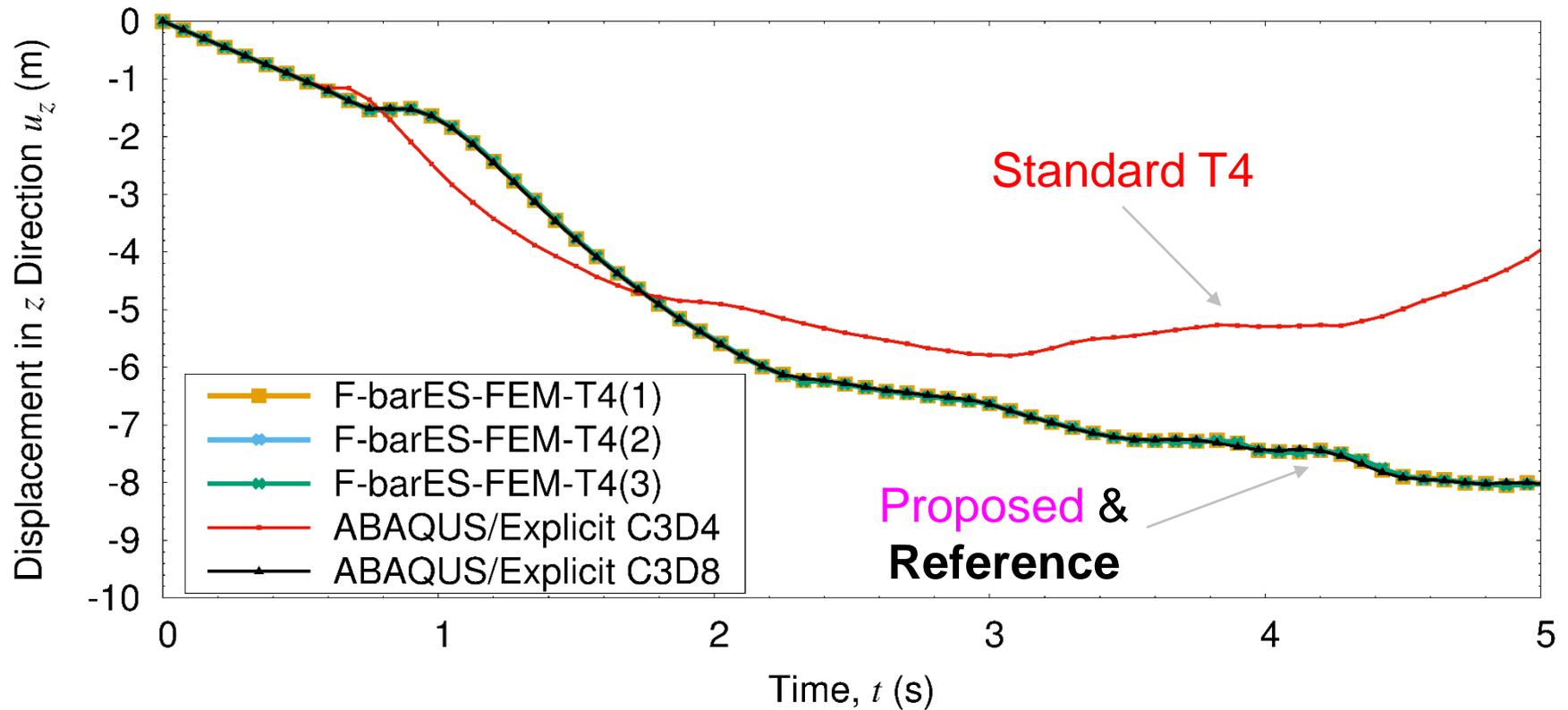
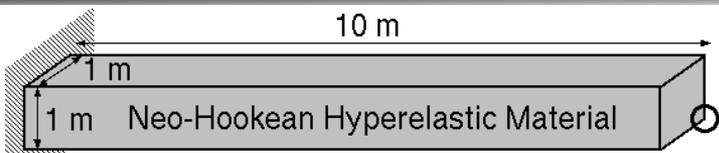
Standard T4 element

✗ Pressure oscillation and locking

F-barES-FEM-T4(2)  
(Proposed method)

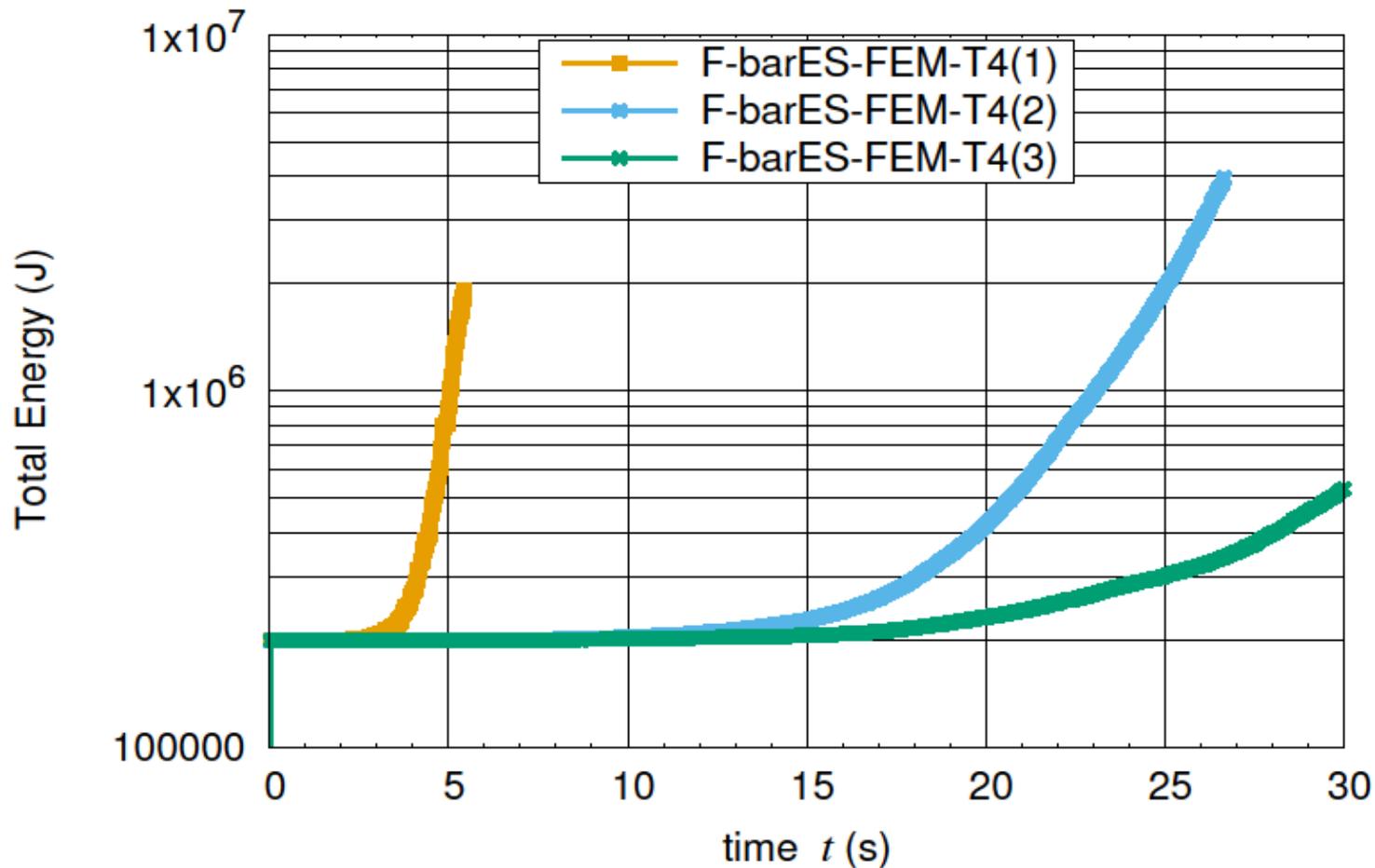
Proposed method shows far better solutions than Standard T4!

# Time history of displacement



- Proposed method shows good result without locking.
- The accuracy of displacement does not depend on the number of cyclic smoothings.

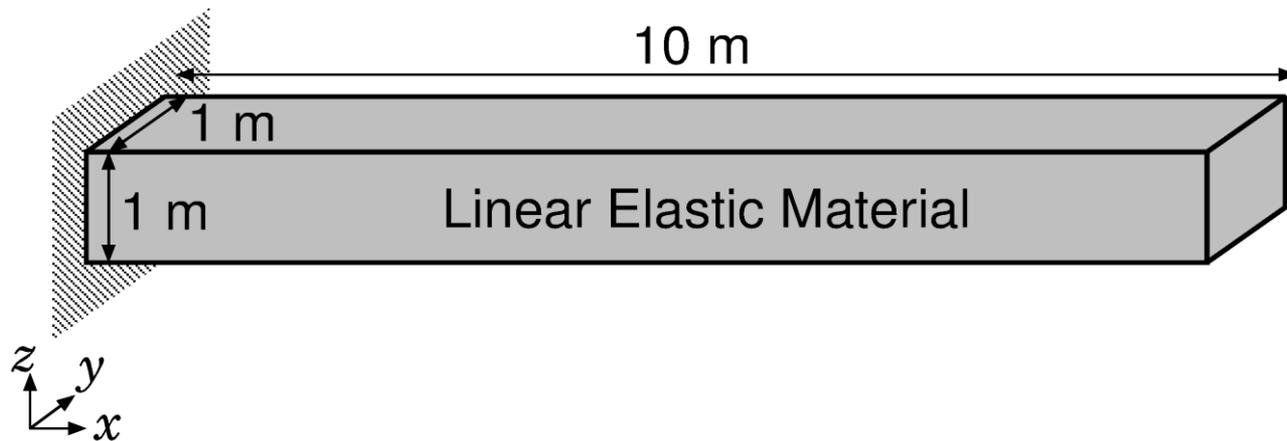
# Time history of total energy



- **Energy divergences** arise in earlier stage...
- Increasing the number of smoothings suppresses the speed of divergence.



# #2 Natural mode of a cantilever



- **Modal** analysis.

- Linear elastic material

Young's modulus: 6.0 MPa,

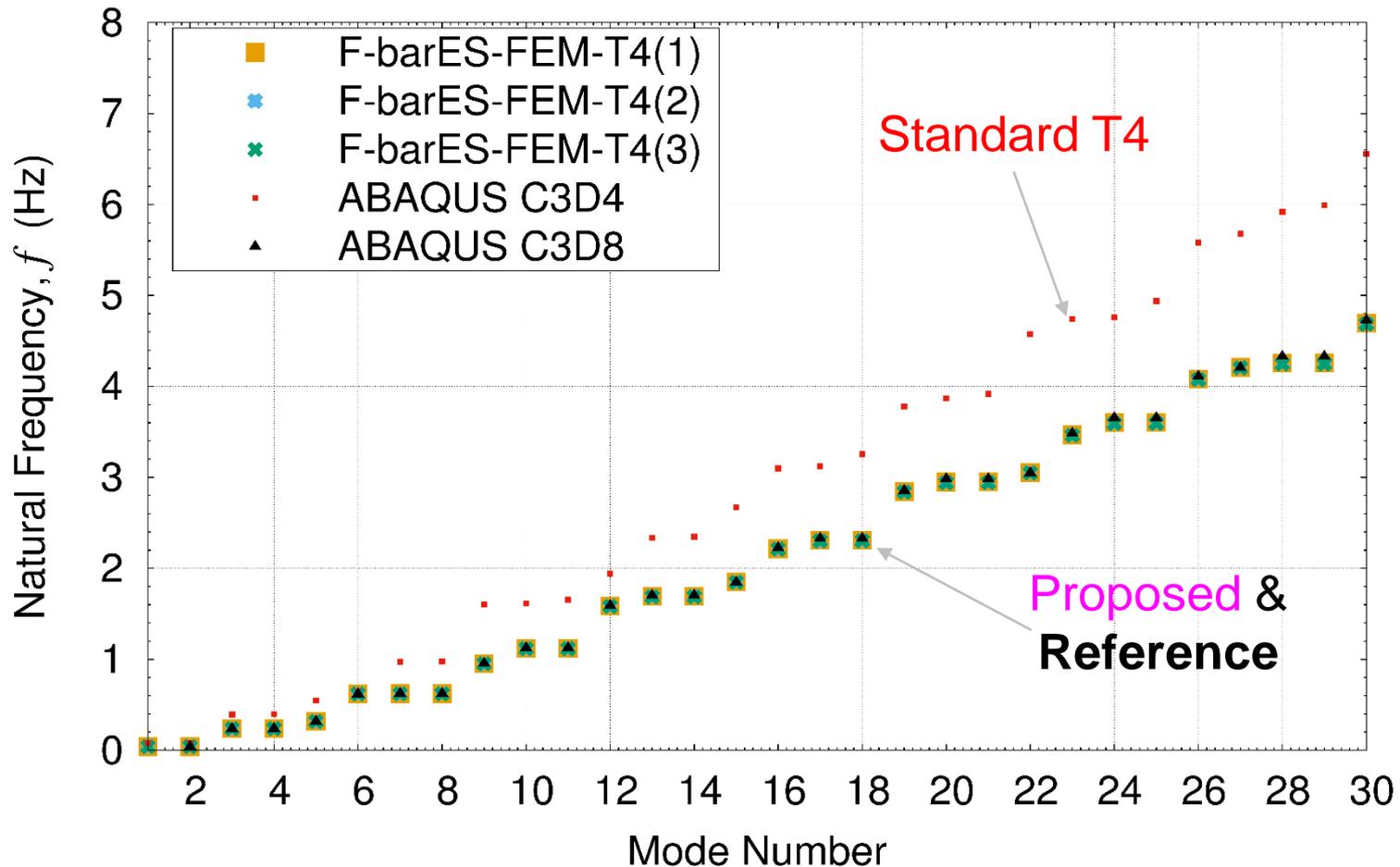
Poisson's ratio: 0.499,

Density: 10000 kg/m<sup>3</sup>.

Same initial elasticity  
as the previous example

- Compare the results of F-barES-FEM-T4, Standard T4 (ABAQUS/Explicit C3D4) and Selective H8 (ABAQUS/Explicit C3D8) elements.

# Natural frequencies of each mode



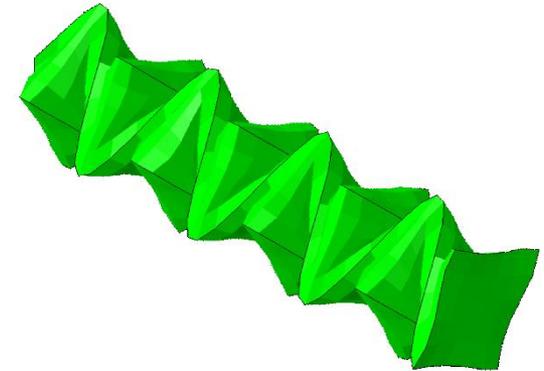
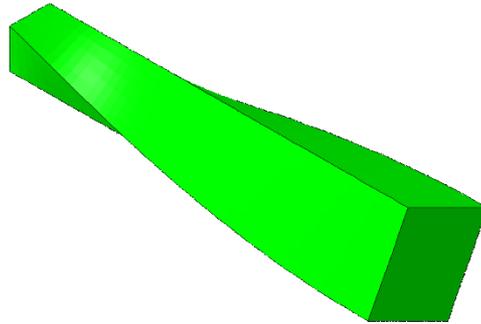
Natural frequencies of our method agree with those of reference !

# Natural mode shapes

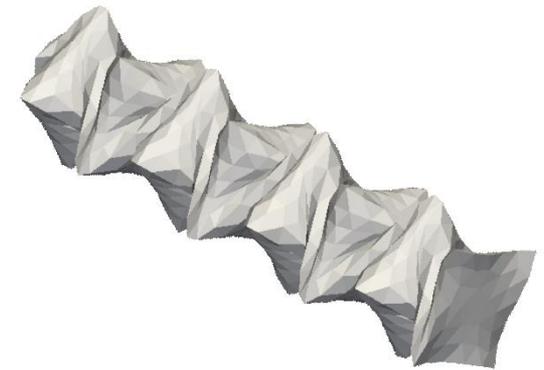
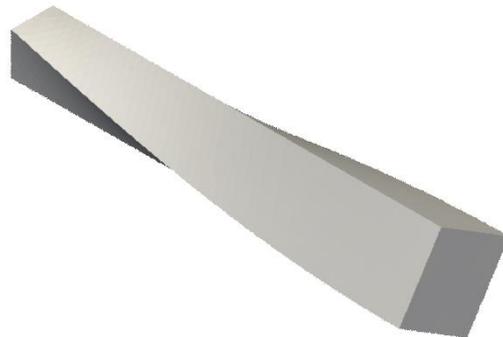
5th

30th

ABAQUS C3D8  
(Reference)

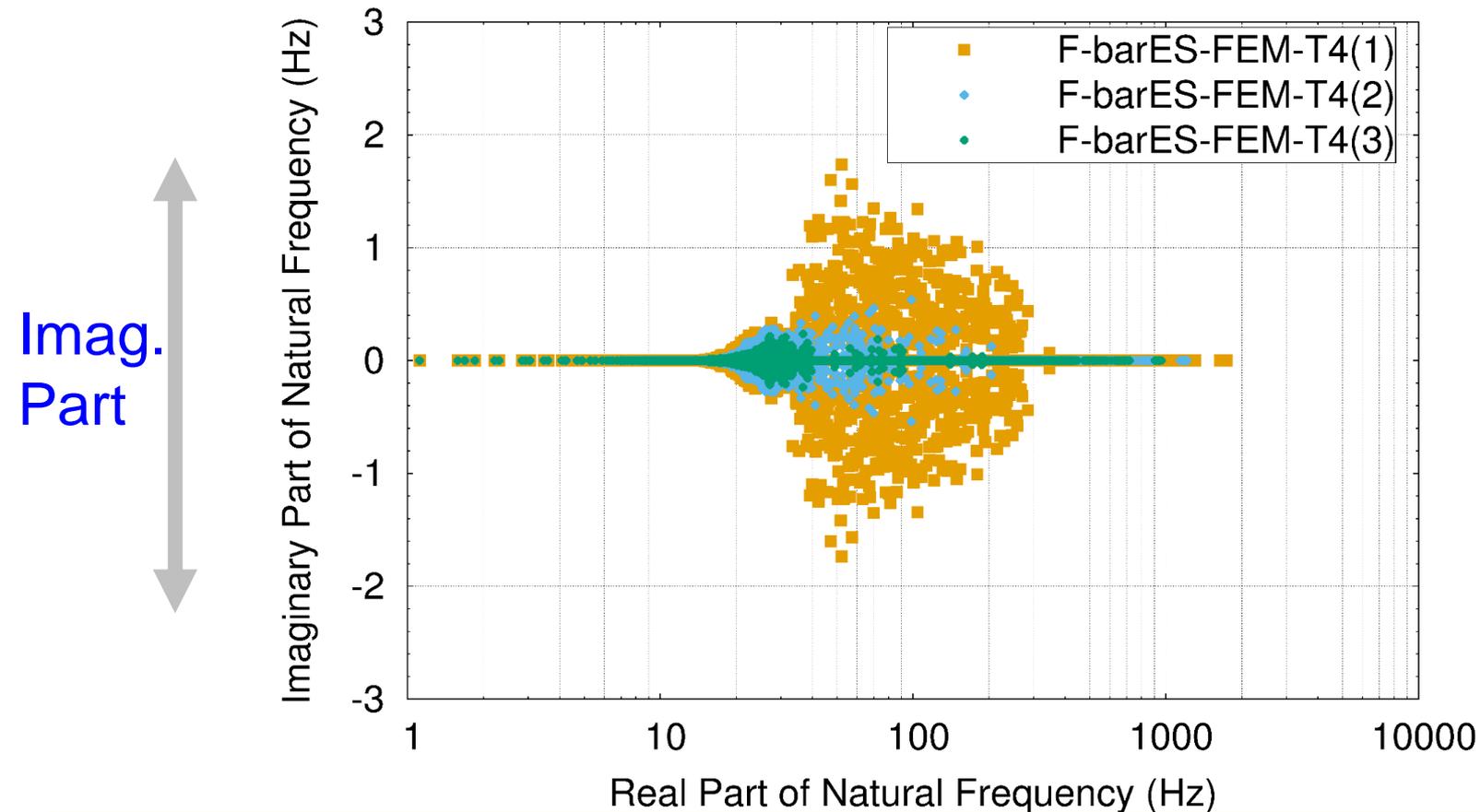


F-barES-FEM(2)  
(Proposed method)



Mode shapes also agree with the reference solutions.

# Distributions of natural frequencies



Some natural frequencies have small **imaginary part**...

Increasing the number of smoothings makes the frequencies close to real numbers.

# Cause of energy divergence

Due to the adoption of F-bar method,  
the stiffness matrix  $[K]$  becomes asymmetric.

Equation of Motion:  $[M]\{\ddot{x}\} + [K]\{x\} = \{f^{\text{ext}}\}$

asymmetric

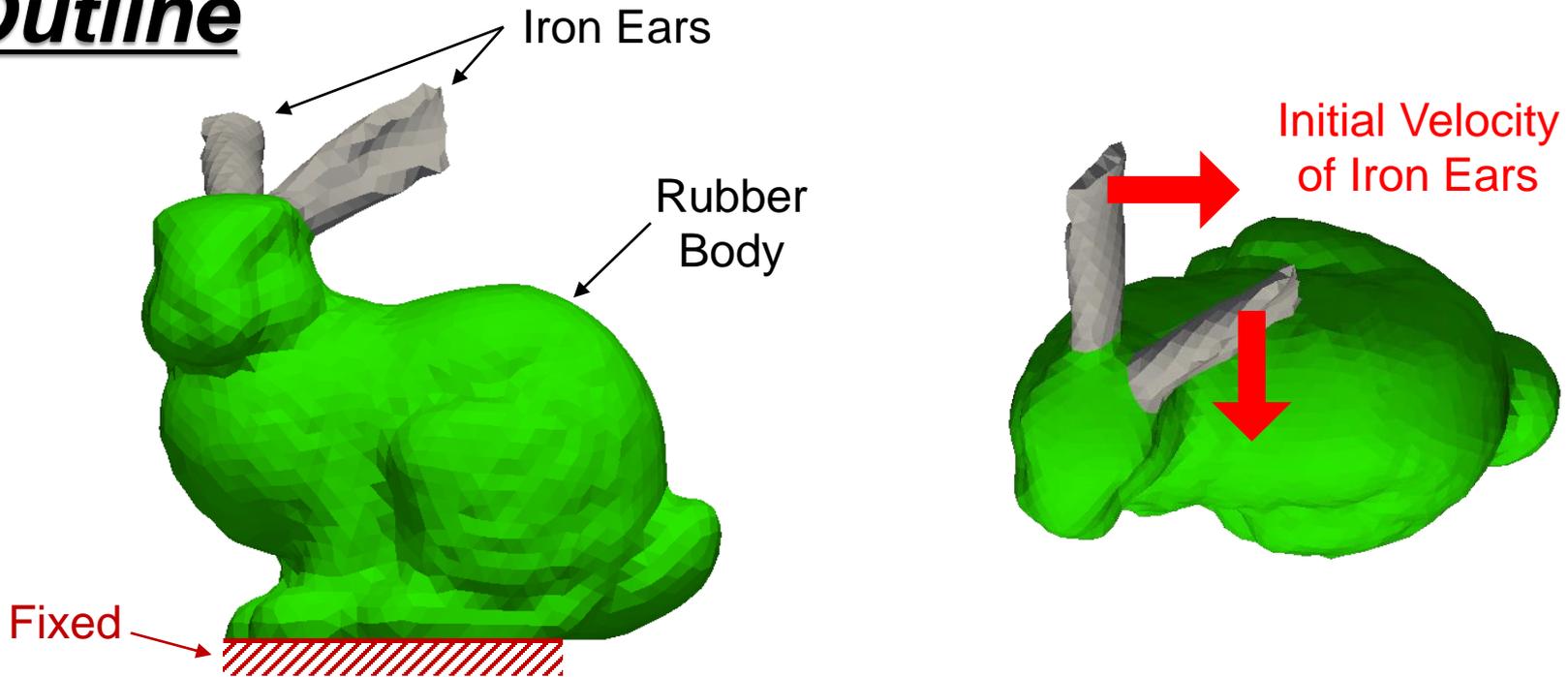
- Asymmetric stiffness matrix gives rise to **imaginary part** of natural frequencies and **instability in dynamic problem**.
- As shown before, increasing the number of smoothings suppress the energy divergence speed.

Our method is restricted to **short-term analysis** (such as impact analysis) with a sufficient number of smoothings.



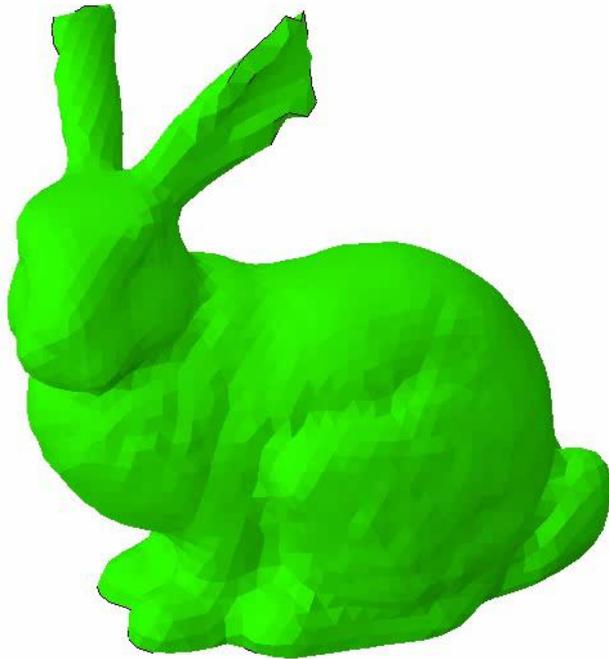
# #3 Swinging of Bunny Ears

## Outline



- Iron ears:  $E_{ini} = 200$  GPa,  $\nu_{ini} = 0.3$ ,  $\rho = 7800$  kg/m<sup>3</sup>, Neo-Hookean, **No cyclic smoothing.**
- Rubber body:  $E_{ini} = 6$  MPa,  $\nu_{ini} = 0.49$ ,  $\rho = 920$  kg/m<sup>3</sup>, Neo-Hookean, **1 cycle of smoothing.**
- Compared to ABAQUS/Explicit C3D4. **No Hex mesh available!**

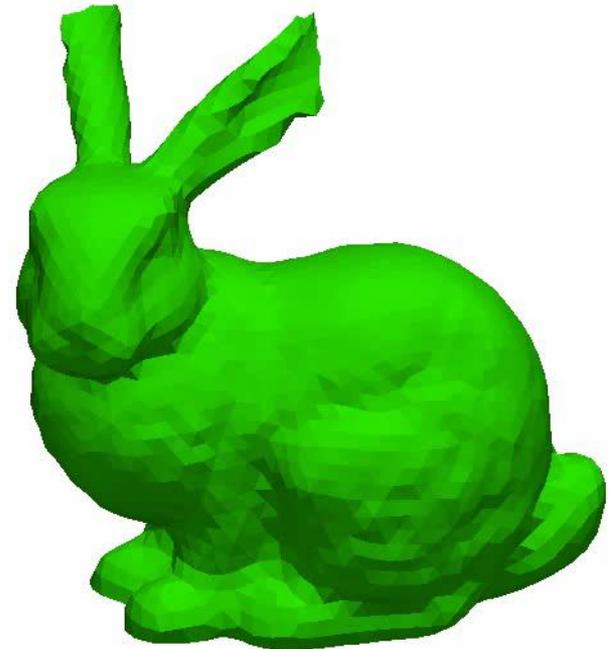
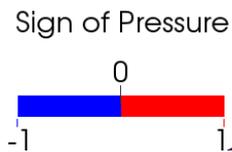
# Time histories of deformed shapes



ABAQUS/Explicit C3D4  
(Standard T4 element)

✗ Pressure oscillation

✗ Locking



F-barES-FEM-T4  
(Proposed method)

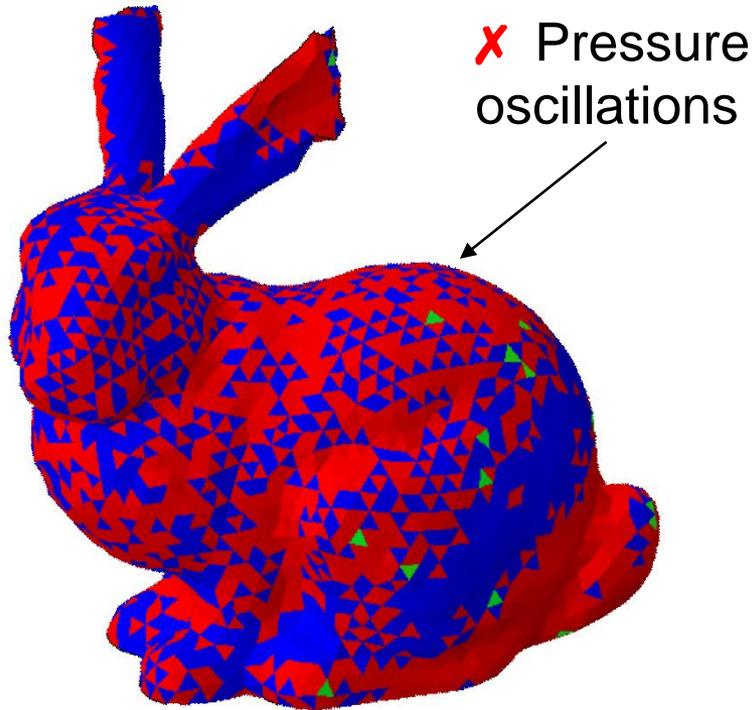
✓ No pressure oscillation

✓ No locking

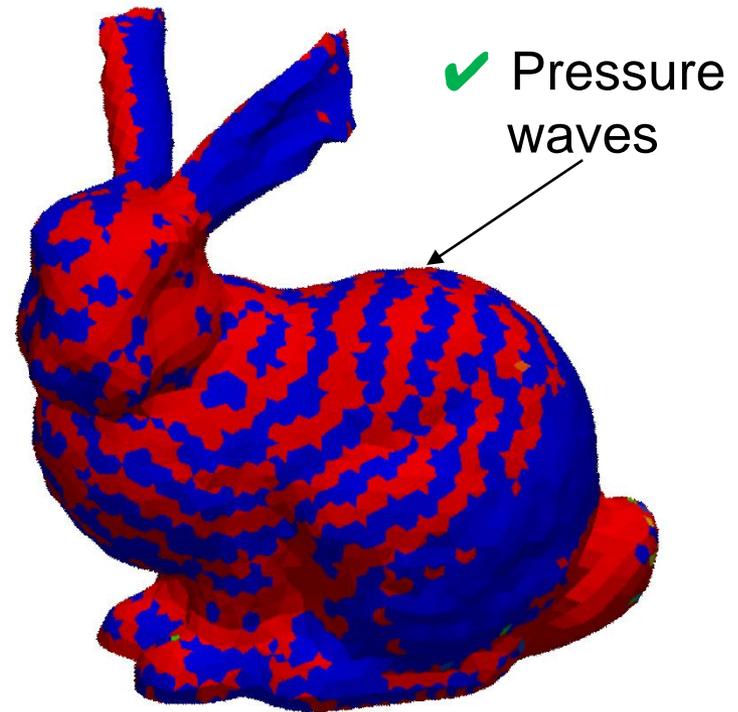
Proposed method seems to be representing not pressure oscillations but pressure waves.

# Deformed shapes and sign of pressure

In an early stage



ABAQUS/Explicit C3D4  
(Standard T4 element)



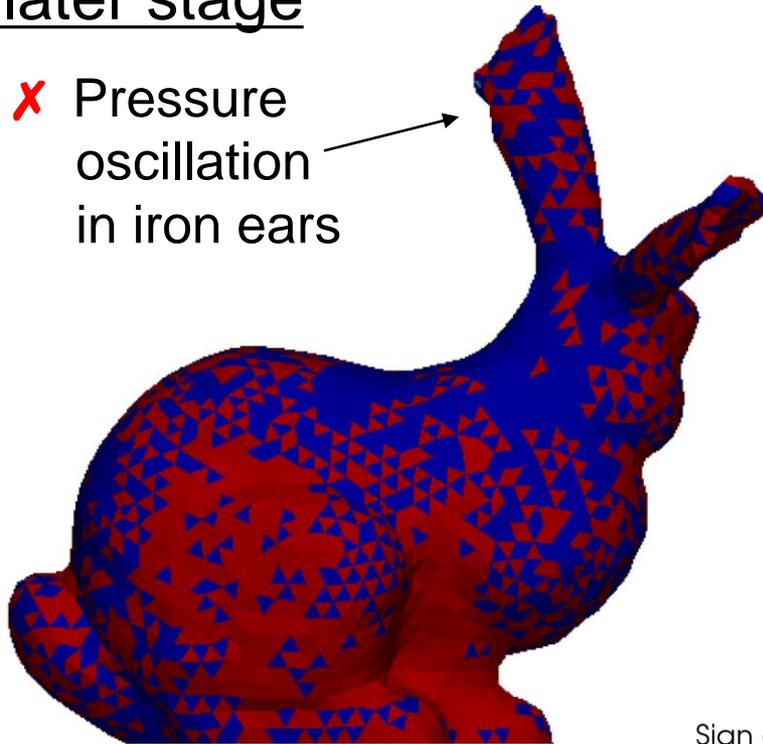
F-barES-FEM-T4  
(Proposed method)

Our method represents pressure waves correctly!

# Deformed shapes and sign of pressure

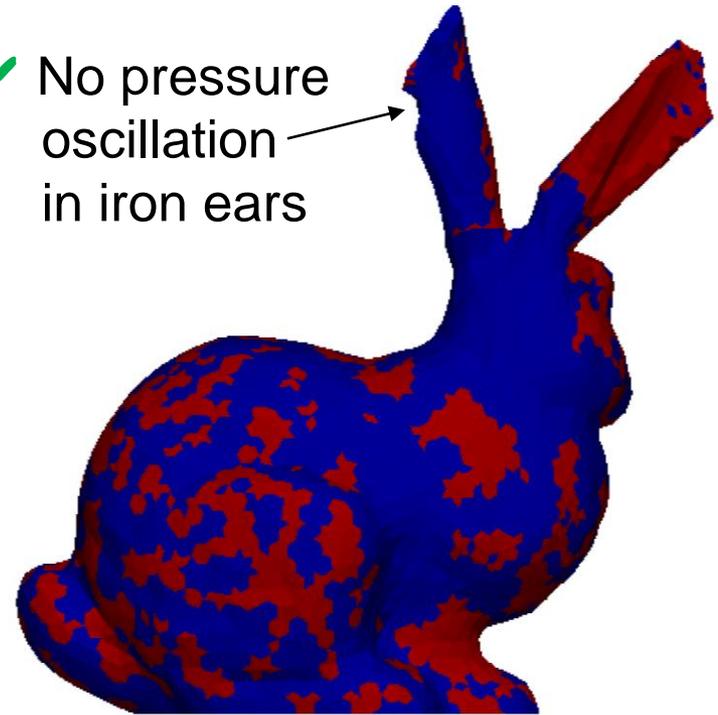
In a later stage

✗ Pressure oscillation in iron ears

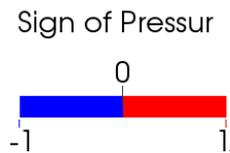


ABAQUS/Explicit C3D4  
(Standard T4 element)

✓ No pressure oscillation in iron ears



F-barES-FEM-T4  
(Proposed method)



It should be noted that a presence of rubber spoils over all accuracy of the analysis with Standard T4 elements.

A rubber parts is a “bad apple” when Standard T4 elements are used.

# Summary

# Summary

- F-barES-FEM-T4 was applied to **dynamic explicit** analysis.
- A few examples of analysis revealed that our method has **excellent accuracy** on relatively **short-term** problems.
- On **long-term problems**, however, our method is **unstable** because of complex natural frequencies.
- Improvement for the long-term stability is our future work.

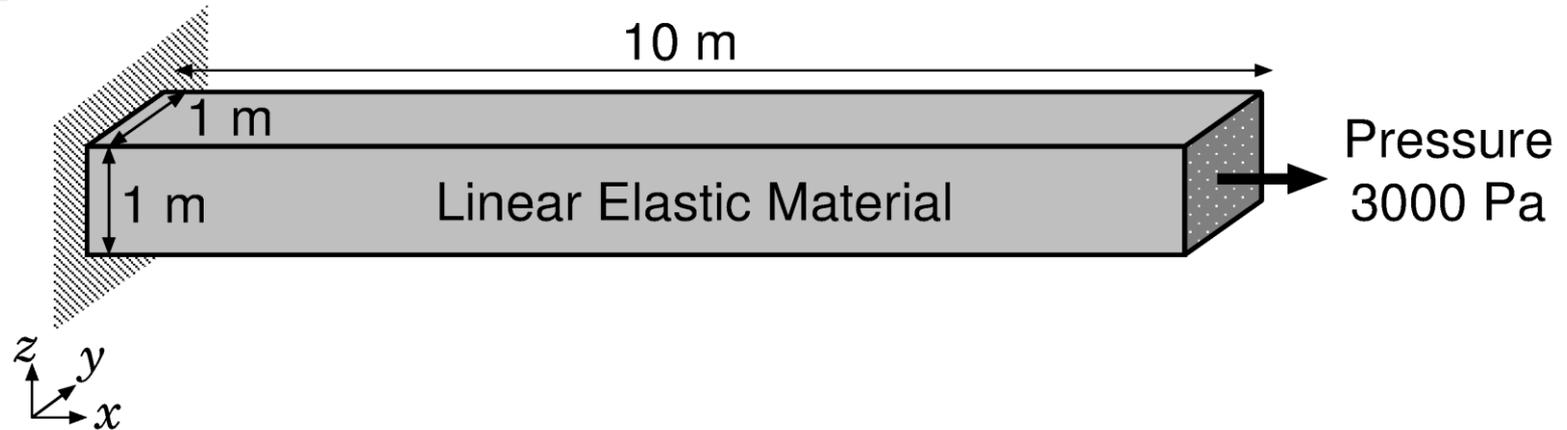
Thank you for your kind attention.

# Appendix



# Propagation of 1D pressure wave

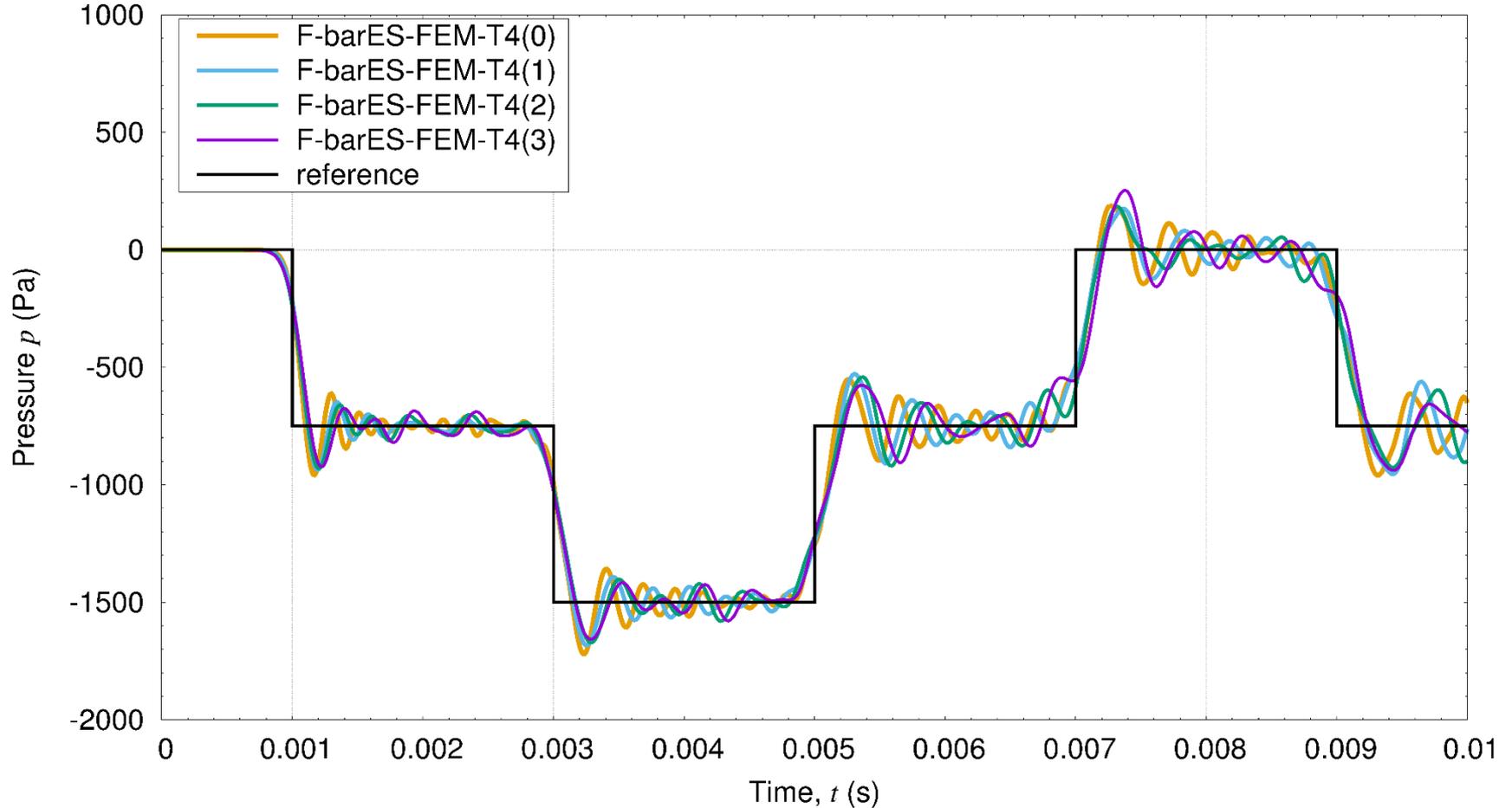
## Outline



- Small deformation analysis.
- Linear elastic material,
  - Young's modulus: 200 GPa,
  - Poisson's ratio: 0.0,
  - Density: 8000 kg/m<sup>3</sup>.
- Results of F-barES-FEM(0), (1), (2), and (3) are compared to the analytical solution.

# Propagation of 1D pressure wave

## Results



# Velocity Verlet Method

## Algorithm

1. Calculate the next displacement  $\{u_{n+1}\}$  as

$$\{u_{n+1}\} = \{u_n\} + \{\dot{u}_n\}\Delta t + \frac{1}{2}\{\ddot{u}_n\}\Delta t^2.$$

2. Calculate the next acceleration  $\{\ddot{u}_{n+1}\}$  as

$$\{\ddot{u}_{n+1}\} = [M^{-1}](\{f^{\text{ext}}\} - \{f^{\text{int}}(u_{n+1})\}).$$

3. Calculate the next velocity  $\{\dot{u}_{n+1}\}$  as

$$\{\dot{u}_{n+1}\} = \{\dot{u}_n\} + \{\ddot{u}_{n+1}\}\Delta t$$

## Characteristics

- 2<sup>nd</sup> order symplectic scheme in time.
- Less energy divergence.

# Cause of energy divergence

Due to the adoption of F-bar method,  
the stiffness matrix  $[K]$  becomes asymmetric  
and thus the dynamic system turns to unstable.

Equation of natural vibration,  $[M]\{\ddot{u}\} + [K]\{u\} = \{0\}$ ,  
derives an eigen equation,  $([M]^{-1}[K])\{u\} = \omega^2\{u\}$ ,  
which has asymmetric left-hand side matrix.

⇒ Some of eigen frequencies could be complex numbers.

⇒ When an angular frequency  $\omega_k = a + ib$  ( $b > 0$ ),  
the time variation of the  $k$ th mode is

$$\begin{aligned}\{u(t)\} &= \text{Re}[\{u_k\} \exp(-i\omega_k t)] \\ &= \text{Re}[\{u_k\} \exp(-iat) \exp(bt)]\end{aligned}$$

Divergent term!